

Study of Dielectric Constant and Loss Tangent of Rochelle Salt Crystal

BALKRISHNA KANDPAL^{1*} and TRILOK CHANDRA UPADHYAY

¹P.D.M. Institute of Engineering & Technology, Karsindhu (Safidon), Jind (Haryana) 126112, India

²Physics Department, H.N.B. Garhwal University (A Central University), Srinagar (Garhwal), Uttarakhand 246174, India

kandpalbk@gmail.com

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Abstract: A two sub-lattice pseudospin-lattice coupled mode model along with third and fourth order phonon anharmonic interactions terms is considered for Rochelle salt. Double Time temperature dependent Green's function method is used for development. Expression for shift, width, renormalized soft mode frequency, Curie temperature, dielectric constant and loss tangent have been derived. Model values of physical quantities are fitted. Theoretical results were compared with experimental results of literature, which show a good agreement.

Keywords: Ferroelectrics, Soft mode, Green function, Anharmonic interactions

Introduction

The Sel de seignette sodium potassium tartrate ($\text{NaKC}_4\text{H}_4\text{O}_6 \cdot 4\text{H}_2\text{O}$) or Rochelle salt is a crystal in which the ferroelectricity was discovered in 1922 by Joseph Valasek. It is colorless to blue while salt crystallizing in the orthorhombic system. It is slightly soluble in alcohol but more completely soluble in water. It has been used as medicinally as a purgative but in more recent years its piezoelectric properties have been important and it has found usage in phonograph pickups and other sensing devices. Its melting point is 75 °C, boiling point 220 °C. Although it is earliest ferroelectric material but is still the subject of intensive study due to its two transitions. It is ferroelectric between 255 K and 297 K showing monoclinic structure in ferroelectric phase. On deuteration the transition temperatures shift to 251 K and 306 K respectively. The theories of ferroelectric properties of Rochelle salt were initiated by Muller¹. After that from microscopic point of view, Mason² assumed that the displacement of the proton in the $\text{O}_{(1)} \cdots (\text{H}_2\text{O})_{(10)}$ hydrogen bond is the ferroelectric dipole and was able to obtain two curie points in agreement with observation. He found it is necessary, however, to assume that the dipole moment itself varies with temperature owing to the thermal expansion. Devonshire³ criticized and modified Mason² theory, but could not reach quite satisfactory results.

Mitsui⁴ proposed a two-sublattice model based on the new information obtained from x-ray studies. Mitsui proposed that ferroelectric dipole moment arises from the ordering of protons in the O --- (H₂O)₍₁₀₎ hydrogen bond. Later on pseudospin model was used by Blinc *et al.*⁵ to explain the two transition temperatures in Rochelle salt. Chaudhuri *et al.*⁶ have used two sublattice-pseudospin lattice coupled model along with a fourth order phonon anharmonic term. Stasyuk and Velychko⁷ have studied dielectric properties of Rochelle salt crystal using a four sublattice model, *i.e.* extension of pseudospin model. However they did not compare their results with experimental data.

Sandy and Jones⁸ have obtained experimental data of dielectric constant and microwave loss tangent for Rochelle salt crystal. Kamba *et al.*⁹ have made experimental studies of vibrational spectroscopy in Rochelle salt. Volkov *et al.*¹⁰ have made experimental study of soft mode frequency in Rochelle salt. Hlinka *et al.*¹¹ have made resonant soft mode study on Rochelle salt. Shiozaki *et al.*¹² have studied structural changes in Rochelle salt due to disordering. Noda *et al.*¹³ have studied calorimetric measurements of phase transition in ammonium mixed Rochelle salt. Kikuta *et al.*¹⁴ have also studied mixed Rochelle salt crystals. Levitskii *et al.*¹⁵ have studied dielectric and piezo, electric properties of Rochelle salt crystal.

Blinc *et al.*⁵ and Kamba *et al.*¹¹ have studied far infrared reflectivity and Raman spectroscopy in Rochelle salt. Volkov *et al.*¹⁰ have made microwave dielectric measurements in Rochelle salt. Hlinka *et al.*¹¹ have made inelastic neutron scattering studies in Rochelle salt crystal.

In present study, a two-sublattice pseudospin-lattice coupled mode model⁹ along with third- and fourth-order phonon anharmonic interactions terms¹⁹ for Rochelle salt has been used. By applying double-time thermal Green's function method²⁰ expressions for shift, width; renormalized soft mode frequency, dielectric constant and loss tangent (microwave absorption) have been evaluated. By using model values of various physical quantities, values of soft mode frequency, dielectric constant and microwave absorption for different temperatures have been calculated and compared with experimental results of Sandy and Jones⁸.

Theory

Following the earlier works^{4,6} the two sublattice pseudospin model with symmetric double-well potential can be expressed as

$$H_s = -2\Omega \sum_i (S_{1i}^x + S_{2i}^x) - \Delta \sum_i (S_{1i}^z + S_{2i}^z) - \sum_{ij} J_{ij} (S_{1i}^z S_{2j}^z + S_{1i}^z + S_{2j}^z) + \sum_{ij} K_{ij} (S_{1i}^z S_{2i}^z), \quad (1a)$$

Where S_i^α is the α th component of the pseudospin $\alpha = x, y, z$ and $l=1$ or 2 corresponding to the different sublattice. The first term of H_s is the tunnelling frequency term with tunnelling frequency Ω and the second term arises due to asymmetry of the crystal field.

To consider the proton-lattice interaction and third and fourth order phonon anharmonic interactions we have to write the total Hamiltonian as

$$H = H_s + H_{s-p} + H_{anh} \quad (1b)$$

$$\text{where } H_{s-p} = -\sum_{ik} V_{ik} (S_{1i}^z A_k + S_i^z A_k^+) \quad (1c)$$

and

$$H_{anh} = \sum_{k_1 k_2 k_3} V^3(k_1, k_2, k_3) A_{k_1} A_{k_2} A_{k_3} + \sum_{k_1 k_2 k_3 k_4} V^4(k_1, k_2, k_3, k_4) A_{k_1} A_{k_2} A_{k_3} A_{k_4} \quad (1d)$$

Following Zubarev¹⁶, we consider the evaluation of Green's function

$$\begin{aligned} G_{ij}(t-t') &= \langle\langle S_{i1}^z(t); S_{j1}^z(t') \rangle\rangle \\ &= -i\theta(t-t') \langle [S_{i1}^z(t); S_{j1}^z(t')] \rangle \end{aligned} \quad (2)$$

Differentiating Eq. (2) twice with respect to time t and t' with the help of Hamiltonian (1b), Fourier transforming it and setting into Dyson's equation form¹⁹.

$$G_{ij}(\omega) = G_{ij}^0(\omega) + G_{ij}^0(\omega) \tilde{P}(\omega) G_{ij}^0(\omega) \quad (3)$$

We obtain Green's function

$$G_{ij}(\omega) = -\frac{\Omega \langle S_{1i}^x \rangle \delta_{ij}}{\pi [\omega^2 - \tilde{\Omega}^2 - 2\Omega i\Gamma(\omega)]} \quad (4)$$

where

$$\hat{\Omega}^2 = \tilde{\Omega}^2 + 2\Omega\Delta(\omega) \quad (5)$$

Eq. (4) is solved self consistently, *i.e.* putting value of $\Delta(\omega)$ into Eq. (4) and

replacing $\omega \rightarrow \tilde{\Omega}$ are gets

$$\hat{\Omega}_{\pm}^2 = \frac{1}{2} \left[(\tilde{\omega}_k^2 + \tilde{\Omega}^2) \pm \sqrt{\{(\tilde{\omega}_k^2 - \tilde{\Omega}^2)^2 + 8V_{ik}^2 \langle S_1^x \rangle \Omega \omega\}} \right] \quad (6)$$

$$\tilde{\Omega}^2 = (a^2 + b^2 - bc) \quad (7)$$

$$a = J_0 \langle S_1^z \rangle + K_0 \langle S_2^z \rangle + \Delta \quad (8)$$

$$b = 2\Omega \quad (9)$$

$$c = 2J_0 \langle S_1^x \rangle + K_0 \langle S_2^x \rangle, \quad (10)$$

$$\tilde{\tilde{\Omega}}^2 = \tilde{\Omega}^2 + 2\Omega\Delta_s(\omega) \quad (11)$$

Dielectric constant ϵ of Rochelle salt as (in ferroelectrics ($\epsilon \gg 1$))

$$\epsilon(\omega) = \frac{-8\pi N\mu^2 \Omega \langle S_{1i}^x \rangle \delta_{ij}}{[\omega^2 - \hat{\Omega}^2] - 2\Omega i\Gamma(\omega)} \quad (12)$$

The tangent loss as

$$\tan \delta = \frac{-2\Omega\Gamma(\omega)}{(\omega^2 - \hat{\Omega}_-^2)} \quad (13)$$

At microwave frequency $\omega(\omega \ll \hat{\Omega}_-)$

$$\tan \delta = \frac{-2\Omega\Gamma(\omega)}{\hat{\Omega}_-^2} \quad (14)$$

Results and Discussion

In order to find temperature dependence of soft mode frequency, dielectric constant and dielectric loss tangent, the model values of physical quantities have been taken from literature⁶. $T_{C1} = 255.2K$, $T_{C2} = 296.9K$, $C_1 = 1830K$, $C_1 = 2248K$, $C_2 = 1830K$, $\eta = 5.5cm^{-1}$, $\Delta = 0.678cm^{-1}$, $\Omega^2(J' + K')^* = 2738cm^{-3}$, $\Omega^2(J' + K') = 2340K$, $\omega_k = 5.20cm^{-1}$, $\Omega V_{ik}^2 = 20.9K$, $A_k K_B \times 10^{17} = 5.73erg/k$, $N = 3.8 \times 10^{21}cm^{-3}$, $\mu = 1.51 \times 10^{18}esu$.

By using these values, $\langle S_i^z \rangle = \frac{J \langle S_i^z \rangle}{\tilde{\Omega}} \tanh \frac{\tilde{\Omega}}{2k_B T}$, $\langle S_i^x \rangle = \frac{\Omega}{\tilde{\Omega}} \tanh \frac{\tilde{\Omega}}{2k_B T}$ have been

calculated for different temperature T (240-330 K) for Rochelle salt crystal. Putting these values in Eq. (7), $\tilde{\Omega}$ is calculated for different temperatures (240-330 K). Now values of $\hat{\Omega}_{\pm}^2$, using Eq. (6), values of ϵ using Eq. (12) and $\tan \delta$ using Eq. (14) have been calculated. Calculated values are shown in Figures 1, 2 and 3 along with data of earlier works⁶ and experimental results of Sandy and Jones⁸.

We have derived expressions for the soft mode frequency, dielectric constant and loss tangent for Rochelle salt crystal. Double time thermal Green's function method and extended two sublattice-pseudospin-lattice coupled mode model along with third-and-fourth-order phonon anharmonic interactions terms have been used in the development.

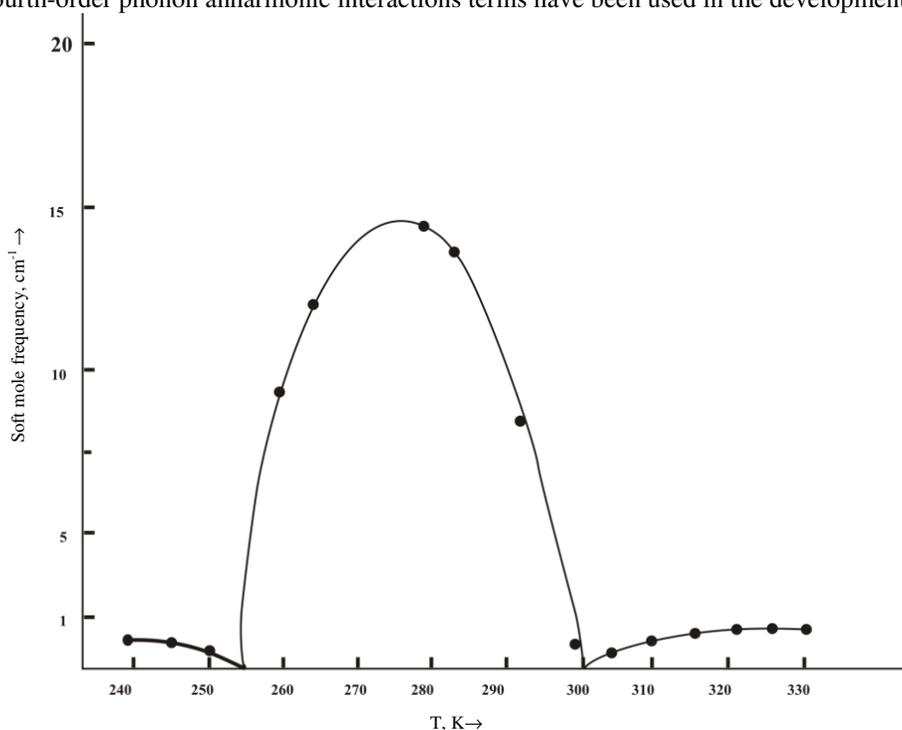


Figure 1. Temperature dependence of soft mode frequency $\hat{\Omega}$ (cm^{-1}) in RS (— our results; • Experimental results of Sandy F and Jones R.V⁸)

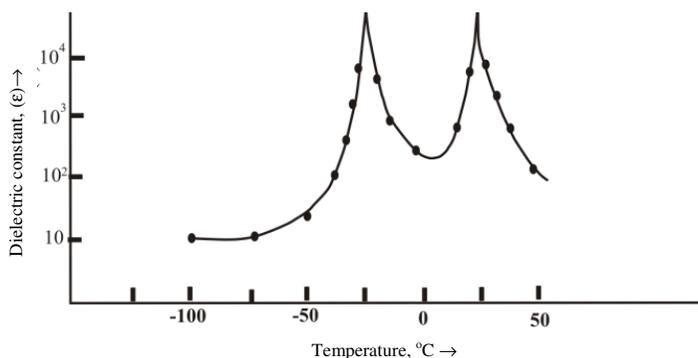


Figure 2. Temperature dependence of dielectric constant (ϵ) in RS (— our results; • Literature value⁸)

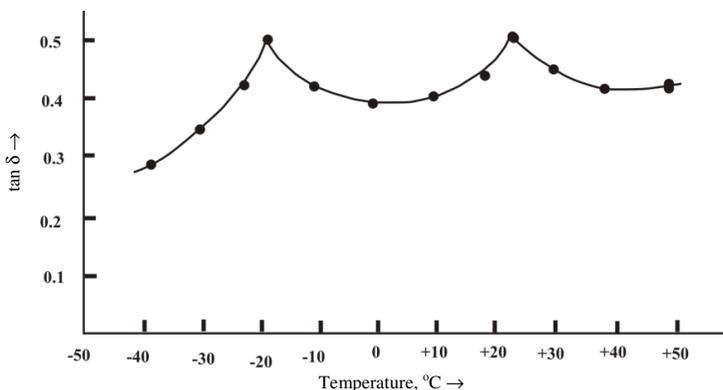


Figure 3. Temperature dependence of loss tangent (δ) in RS (— our results; • Literature value⁸)

Model values from literature are fitted in the expressions and temperature dependence of soft mode, dielectric constant and loss tangent are calculated. Theoretical results are compared with theoretical results of earlier workers⁶ and experimental data of Sandy and Jones⁸. Our procedure differs earlier works in the sense that they have not considered third-order phonon anharmonic interactions. They decoupled the correlations in the early beginning. The width and shift and the third-order phonon interactions are contributions of present work. If the width, shift and third order phonon anharmonic interactions terms are neglected from present calculations, our results reduce to the results of Chaudhari *et al*⁶. Our expressions predicts that loss first increases from below, upto transition then decreases and again increases upto second transition and then decreases. A transverse radiation field derives the low laying transverse mode of the material in a forced vibration. Energy is transferred from the electromagnetic field to this lattice mode and is then degraded into other vibrational modes of the material. Due to anharmonic phonon interactions, decay processes take place. For example, third-order interaction leads to the decay of a virtual phonon into two real phonons or the virtual phonon may destroyed by scattering a thermally excited phonon. Similar processes occur for fourth-order interactions.

Conclusion

As can be understood from the above results, the pseudospin-lattice coupled mode model along with third-and fourth- order phonon anharmonic interactions terms explains the dielectric

properties of Rochelle salt clearly. The phonon anharmonic interaction terms significantly affect the temperature dependence of soft mode frequency, dielectric constant and loss tangent in Rochelle salt. These results are better than results of earlier authors, quantitatively.

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References

1. Muller K A, *Helv Phys Acta*, 1986, **59**, 847.
2. Mason W P, *Phys Rev.*, 1947, **72**, 854; DOI:10.1103/PhysRev.72.854
3. Devonshire A F, *Phil Mag.*, 1957, **2(20)**, 1027-1039; DOI:10.1080/14786435708238209
4. Mitsui T, *Phys Rev.*, 1958, **111**, 1259-1267; DOI:10.1103/PhysRev.111.1259.
5. Blinc R, Petkovsek J and Zupancic I, *Phys Rev.*, 1964, **136**, A1684; DOI:10.1103/PhysRev.136.A1684
6. Chaudhari B K, Atake T, Ganguli S and Chihara H, *J Phys Soc Jpn.*, 1980, **49**, 608; DOI:10.1143/JPSJ.49.608
7. Stasyuk I V and Velychko O V, *J Phys C*, 2004, **16**, 1963.
8. Sandy F and Jones R V, *Phys Rev.*, 1968, **168**, 481; DOI:10.1103/PhysRev.168.481
9. Kamba S, Schaak G and Petzelt J, *Phys Rev B*, 1995, **51**, 14998; DOI:10.1103/PhysRevB.51.14998
10. Volkov A A, Kozlov G V, Kryukova E B and Petzelt J, *Zh Eksp Teor Fiz*. 1986, **90(1)**,192–200. http://www.jetp.ac.ru/cgi-bin/dn/e_063_01_0110.pdf
11. Hlinka J, Kulda J, Kamba S and Petzelt J, *Phys Rev B*, 2001, **63**, 052102; DOI:10.1103/PhysRevB.63.052102
12. Shiozaki Y, Shimizu K and Nozaki R, *Ferroelctrics*, 2001, **261(1)**, 239-244; DOI:10.1080/00150190108216291
13. Noda N, Nozaki R and Shiozaki Y, *Phys Rev B*, 2000, **62**, 12040; DOI:10.1103/PhysRevB.62.12040
14. Kikuta T, Kawabe R, Yamazaki T and Nakatami N, *J Korean Phys Soc.*, 2003, **42**, 51275.
15. Levitskii R R, Zukek I R, Verkholyak T M and Moina A P, *Phys Rev B*, 2003, **67** 174112; DOI:10.1103/PhysRevB.67.174112
16. Zubarev D N, *Sov Phys Usp.*, 1960, **3**, 320; DOI:10.1070/PU1960v003n03ABEH003275
17. Semwal B S and Sharma P K, *Prog Theor Phys Jpn.*, 1974, **51(3)**, 639-655; DOI:10.1143/PTP.51.639
18. Born M and Haung K, *Dynamical Theory of Crystal Lattices*, Oxford Press New York, 1954.
19. Upadhyay T C and Semwal B S, *Pramana J Phys.*, 2003, **60(3)**, 525-533; DOI:10.1007/BF02706161
20. Upadhyay T C, Bhandari R S and Semwal B S, *Pramana J Phys.*, 2006, **67(3)**, 547-552; DOI:10.1007/s12043-006-0016-y
21. Kubo R, *J Phys Soc Jpn.*, 1957, **12**, 570-586; DOI:10.1143/JPSJ.12.570